



# Application of solid-state NMR restraint potentials in membrane protein modeling

Jinhyuk Lee<sup>a</sup>, Jianhan Chen<sup>b</sup>, Charles L. Brooks III<sup>c</sup>, Wonpil Im<sup>a,\*</sup>

<sup>a</sup>Department of Molecular Biosciences and Center for Bioinformatics, The University of Kansas, 2030 Becker Drive, Lawrence, KS 66047, USA

<sup>b</sup>Department of Biochemistry, The Kansas State University, 141 Chalmers Hall, Manhattan, KS 66506, USA

<sup>c</sup>Department of Chemistry, The University of Michigan, 930 North University Avenue, Ann Arbor, MI 48109, USA

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## ABSTRACT

We have developed a set of orientational restraint potentials for solid-state NMR observables including <sup>15</sup>N chemical shift and <sup>15</sup>N–<sup>1</sup>H dipolar coupling. Torsion angle molecular dynamics simulations with available experimental <sup>15</sup>N chemical shift and <sup>15</sup>N–<sup>1</sup>H dipolar coupling as target values have been performed to determine orientational information of four membrane proteins and to model the structures of some of these systems in oligomer states. The results suggest that incorporation of the orientational restraint potentials into molecular dynamics provides an efficient means to the determination of structures that optimally satisfy the experimental observables without an extensive geometrical search.

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## 1. Introduction

Helices are the major secondary structures in transmembrane (TM) domains of membrane proteins [1], which play important roles in intracellular and intercellular signaling, transportation of ions and small molecules, and cell–cell recognition [2,3]. Identifying and characterizing the relative orientation of helices in membrane proteins are crucial in determining their topologies and three-dimensional structures, which can provide insights into the underlying pathologic mechanisms responsible for many human diseases. Although particular successes are evident for some membrane proteins [4–6], X-ray diffraction methods generally encounter numerous obstacles due to hydrophobic interactions between membrane proteins and lipids. Solid-state NMR (SSNMR) is an emerging technique to study the topology, structure, and dynamics of membrane proteins in their native environment of lipid bilayers, which can complement X-ray and solution NMR studies. Measurements of <sup>15</sup>N chemical shift and <sup>15</sup>N–<sup>1</sup>H dipolar coupling using various SSNMR techniques provide the orientational constraints of helices that can be used to determine their tilt and rotation angles with respect to lipid bilayers. Such orientational information can also be utilized to model assembled structures in certain oligomer states.

Monte Carlo (MC) approach based on a geometric search of tilt and rotation angles have been widely used to find global minimum structures in the energy (or penalty) function that incorporates all the orientational constraints and the empirical potential energy function [7–9]. Alternatively, the incorporation of the orientational restraint potentials into molecular dynamics (MD) simulations has been attempted and applied in the structure refinement of gramicidin A [10]. Although the use of restraint potentials in energy minimization and MD simulations has been well-established as a powerful tool in the solution NMR structure determination, similar approaches have not been fully explored in the context of the SSNMR structure determination.

The present work aims to develop a protocol to determine membrane protein structures using MD simulations with orientational restraint potentials representing SSNMR observables such as the chemical shift and dipolar coupling. For this, we first formulated and implemented these restraint potentials as well as corresponding analytical forces in the biomolecular simulation program CHARMM [11]. Then, we developed and optimized a protocol for SSNMR structure determination by applying the restraint potentials to several membrane proteins whose structures were previously determined using the MC approach. To illustrate the efficacy of our approach, we determined the monomer structures of the major pVIII coat protein of fd filamentous bacteriophage (fd coat) [12], the mercury ion transporter (MerF) from the bacterial mercury detoxification system [13], the TM domain of the M2 protein (M2TMP) from influenza A virus [14], and the TM domain

\* Corresponding author. Fax: +1 785 864 5558.

E-mail address: [wonpil@ku.edu](mailto:wonpil@ku.edu) (W. Im).

of viral protein “u” (VpuTM) from human immunodeficiency virus 1 (HIV-1) [9]. We further utilized the monomer structures to model oligomeric states of VpuTM (tetramer and pentamer) and M2TMP (tetramer).

## 2. Development of SSNMR restraint potentials

### 2.1. Dipolar coupling

Dipolar coupling results from the interaction of one nuclear spin  $i$  with a magnetic field generated by another nuclear spin  $j$ , or vice versa. This is an interaction through space that is dependent on the gyromagnetic ratio ( $\gamma$ ) of each nucleus as well as the distance,  $r_{ij}$ . The individual  $^{15}\text{N}$ - $^1\text{H}$  dipolar coupling ( $v$ ) is measured as

$$v = \frac{v_0}{2} (3 \cos^2 \theta - 1), \quad (1)$$

where  $v_0$  is the dipolar coupling constant, i.e.,  $v_0 = (\gamma_{\text{N}}\gamma_{\text{H}}h\mu_0)/(8\pi^3r_{\text{NH}}^3)$  in which  $h$  is Planck's constant and  $\mu_0$  is the permeability of space. In the case of  $r_{\text{NH}} = 1.07 \text{ \AA}$ , the dipolar coupling constant becomes 19.86 kHz.  $\theta$  is the angle between the NH internuclear vector,  $\mathbf{r}_{\text{NH}}$ , and the magnetic field, which is generally assumed to be parallel to the Z-axis. The solution of Eq. (1) has an intrinsic degeneracy due to doublet splitting of the dipolar coupling. Furthermore, SSNMR experiments cannot distinguish the sign of observables below half maximal amplitude, i.e.,  $v_0/2$ , which is the so-called sign degeneracy [10,15]. In spite of these ambiguities, dipolar waves have been used to identify a helix in proteins, detect kinks or curvature in the helix, and determine rotation and orientation of the helix with respect to the membrane normal [16].

The restraint potential for  $^{15}\text{N}$ - $^1\text{H}$  dipolar coupling,  $U_{\text{dc}}$ , is expressed as a quadratic function to restrain the deviation of the calculated values ( $v$ ) from the experimental observables ( $v_{\text{exp}}$ ):

$$U_{\text{dc}} = \sum_i^N k_{\text{dc}} (|v_i| - v_{\text{exp},i})^2, \quad (2)$$

where  $N$  is the number of dipolar couplings, and  $k_{\text{dc}}$  is the force constant in kcal/(mol kHz<sup>2</sup>). Note that an absolute value,  $|v|$ , is used in Eq. (2) to account for the sign degeneracy below  $v_0/2$ . The derivatives of  $U_{\text{dc}}$  with respect to  $\mathbf{r}_\alpha$  of two selected atoms ( $\alpha = \text{N}$  or  $\text{H}$ ) become

$$\frac{\partial U_{\text{dc}}}{\partial \mathbf{r}_\alpha} = \begin{cases} 2 \sum_i^N k_{\text{dc}} (v_i - v_{\text{exp},i}) \frac{\partial v_i}{\partial \mathbf{r}_\alpha} & \text{if } 0 \leq v \leq v_0 \\ 2 \sum_i^N k_{\text{dc}} (v_i + v_{\text{exp},i}) \frac{\partial v_i}{\partial \mathbf{r}_\alpha} & \text{if } -\frac{v_0}{2} \leq v \leq 0. \end{cases} \quad (3)$$

The derivative of  $|v_i|$  in Eq. (2) yields an opposite sign when  $v_i < 0$ . If the sign degeneracy is resolved, the derivatives can be described by the upper equation in Eq. (3). The last term,  $\partial v_i / \partial \mathbf{r}_\alpha$ , becomes

$$\frac{\partial v_i}{\partial \mathbf{r}_\alpha} = 3v_{0,i} \cos \theta_i \frac{\partial \cos \theta_i}{\partial \mathbf{r}_\alpha} + \frac{3 \cos^2 \theta_i - 1}{2} \frac{\partial v_{0,i}}{\partial \mathbf{r}_\alpha}, \quad (4)$$

where  $\partial v_{0,i} / \partial \mathbf{r}$  does not vanish when the NH distance varies during simulations. The detailed mathematical descriptions for the derivatives of  $\cos \theta$  and  $v_0$  are given in Supplementary Information (SI).

### 2.2. SSNMR chemical shift

Chemical shielding is an anisotropic interaction characterized by a shielding tensor, which can be diagonalized to yield three principal components ( $\sigma_{11}$ ,  $\sigma_{22}$ , and  $\sigma_{33}$ ). The three components in the  $^{15}\text{N}$  chemical shift are described in terms of a peptide plane (molecular frame) defined by the N, H, and C atoms. The shielding or de-shielding effects from the peptide plane move a resonant frequency to either higher or lower frequencies from a reference

frequency. This deviation from the reference frequency is known as chemical shift. To calculate the chemical shift efficiently, the chemical shift tensor needs to be fixed relative to a local molecular frame, and the magnitude of each component must be constant, which is the so-called “rigid tensor approximation” [17,18]. Because the SSNMR time scale is much longer than that of molecular motions, the observed chemical shift ( $\sigma_{\text{exp}}$ ) parallel to the magnetic field can be calculated as a time-average of the projected instantaneous second-rank tensor,

$$\sigma_{\text{exp}} = \left\langle \sum_{i=1}^3 [\hat{n} \cdot \hat{e}_i(t)] \sigma_{ii}(t) [\hat{e}_i(t) \cdot \hat{n}] \right\rangle, \quad (5)$$

where  $\hat{n}$  is a unit vector of the magnetic field, which is assumed to be parallel to the membrane normal (the Z-axis;  $\hat{n} \equiv \hat{z}$ ), and  $\sigma_{ii}(t)$  and  $\hat{e}_i(t)$  are the instantaneous magnitude and unit vector of chemical shift tensors, respectively. In general,  $\hat{e}_1$  and  $\hat{e}_3$  are on the peptide plane defined by N, C, and H atoms, and  $\hat{e}_2$  is defined by the cross product of  $\mathbf{r}_{\text{NC}}$  and  $\mathbf{r}_{\text{NH}}$  [18]. Then,  $\hat{e}_1$  is defined by a rotation angle  $\phi$  from  $\mathbf{r}_{\text{NH}}$  on the peptide plane, and  $\hat{e}_3 = \hat{e}_1 \times \hat{e}_2$ . In terms of the rigid tensor approximation, the magnitudes of chemical shift tensor are separated from their unit vectors, i.e.,

$$\sigma_{\text{exp}} = \sum_{i=1}^3 \langle [\hat{z} \cdot \hat{e}_i(t)]^2 \rangle \sigma_{ii}. \quad (6)$$

The experimentally observed chemical shift is the time-average of the individual chemical shift tensors of the peptide plane. However, the chemical shift ( $\sigma$ ) of a molecule can be calculated from an instantaneous conformation in the course of its time evolution. Then, the chemical shift parallel to the magnetic field becomes

$$\sigma = \sigma_{11} \hat{e}_{1,z}^2 + \sigma_{22} \hat{e}_{2,z}^2 + \sigma_{33} \hat{e}_{3,z}^2, \quad (7)$$

where  $\hat{e}_{i,z}$  is the  $z$  component of  $\hat{e}_i$  ( $i = 1, 2, 3$ ). The magnitudes of the three chemical shift tensors ( $\sigma_{11}$ ,  $\sigma_{22}$ , and  $\sigma_{33}$ ) are often assumed to be independent of amino acid residues, except for glycine and proline. Table 1 summarizes these parameters, including  $\phi$  to define  $\hat{e}_1$  in our testcases.

The chemical shift restraint potential is expressed as a quadratic function to restrain deviations of calculated chemical shifts ( $\sigma$ ) from experimental observables ( $\sigma_{\text{exp}}$ );

$$U_{\text{cs}} = \sum_i^N k_{\text{cs}} (\sigma_i - \sigma_{\text{exp},i})^2, \quad (8)$$

**Table 1**  
Summary of SSNMR parameters for four testcases

	Fd coat	M2TMP	VpuTM	MerF
No. of chemical shifts <sup>a</sup>	38	10	18	41 <sup>b</sup>
No. of dipolar couplings <sup>a</sup>	39	7	18	41 <sup>b</sup>
No. of residues	50	25	19	45
PDB (Monomer)	1MZT	1MP6	1PJE	2H30
PDB (Oligomer)		1NYJ	1PI8, 1PI7	
References	[12]	[7,41]	[9]	[9,13]
$\sigma_{11}$ (ppm)	64/41 <sup>c</sup>	Indiv. <sup>d</sup>	64	64/41/38 <sup>e</sup>
$\sigma_{22}$ (ppm)	77/64	Indiv.	77	77/64/127
$\sigma_{33}$ (ppm)	217/210	Indiv.	222	222/215/231
$\phi$ (°)	107/108	109	108.5	108.5/108.5/110
Dist(NH) (Å)	1.07/1.07	1.07	1.05	1.05/1.05/NA <sup>f</sup>

<sup>a</sup> The SSNMR observables are obtained from corresponding references.

<sup>b</sup> The SSNMR observables are obtained in magnetically aligned bicelles. The corresponding observables in bilayers can be calculated by  $\sigma_{\text{bicelle}} = (1 + S/2)\sigma_{\text{iso}} - (S/2)\sigma_{\text{bilayer}}$  and  $v_{\text{bicelle}} = -(S/2)v_{\text{bilayer}}$ , where  $\sigma_{\text{iso}}$  is the isotropic chemical shift frequency. The order parameter  $S$  was set to 0.80 for the bicelles.

<sup>c</sup> The chemical shift tensors for non-glycine/glycine residues.

<sup>d</sup> Different chemical shift tensors were used for individual residues. The values are in the references.

<sup>e</sup> The chemical shift tensors for non-glycine/glycine/proline residues.

<sup>f</sup> There is no NH bond in the Pro backbone.

where  $N$  is the number of chemical shifts, and  $k_{cs}$  is the force constant in kcal/(mol ppm<sup>2</sup>). The derivatives of the chemical shift restraint potential ( $U_{cs}$ ) with respect to  $\mathbf{r}_\alpha$  ( $\alpha = \text{N, H, or C}$ ) are

$$\frac{\partial U_{cs}}{\partial \mathbf{r}_\alpha} = 2 \sum_i^N k_{cs} (\sigma_i - \sigma_{\text{exp},i}) \frac{\partial \sigma_i}{\partial \mathbf{r}_\alpha}. \quad (9)$$

In terms of Eq. (7), the last term for  $i$ th chemical shift becomes

$$\frac{\partial \sigma}{\partial \mathbf{r}_\alpha} = 2 \left( \sigma_{11} \hat{e}_{1,z} \frac{\partial \hat{e}_{1,z}}{\partial \mathbf{r}_\alpha} + \sigma_{22} \hat{e}_{2,z} \frac{\partial \hat{e}_{2,z}}{\partial \mathbf{r}_\alpha} + \sigma_{33} \hat{e}_{3,z} \frac{\partial \hat{e}_{3,z}}{\partial \mathbf{r}_\alpha} \right), \quad (10)$$

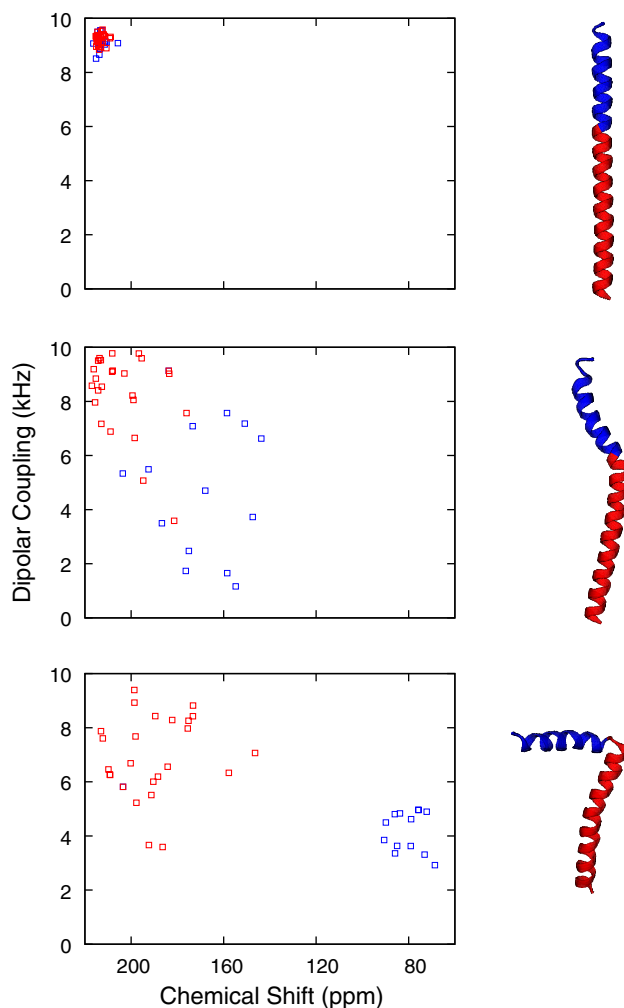
where the index  $i$  is omitted for the sake of clarity. The detailed mathematical description of each derivative is described in SI.

### 3. Simulation details

To examine the numerical accuracy of the implemented analytical forces, we first compared them with the finite difference forces. Both forces are consistent within numerical errors (data not shown), demonstrating that the analytical forces are implemented correctly. In general, the simulation stability with restraint potentials is quantified by measuring the so-called total energy fluctuation in NVE (constant volume and energy) simulations [19,20]. We performed 200 ps NVE simulations of VpuTM with the Verlet integrator and a time step of 1 femtosecond (fs). As shown in SI Fig. S1, the total energies appear to be well conserved, demonstrating again the accuracy of calculated restraint forces. SI Figs. S2 and S3 show the total energy fluctuations from NVE Cartesian MD and torsion angle MD (TAMD) simulations [19] as a function of restraint force constant. As expected, the simulation stability depends on the applied force constants and dynamics algorithms. TAMD appears to yield much more stable trajectories than Cartesian MD because the peptide plane is a rigid unit in TAMD.

Combined with constant temperature TAMD simulations at 300 K, we have developed and optimized a protocol to determine monomer structures using the SSNMR restraint potentials. The protocol consists of (1) generating an ideal  $\alpha$ -helix (or helices), (2) performing several steps of TAMD simulations with increasing force constants, and (3) selecting structures that satisfy certain criteria such as  $\delta\sigma \leq 1$  ppm and  $\delta\nu \leq 1$  kHz, where  $\delta\sigma$  and  $\delta\nu$  represent the root-mean square deviations of the calculated chemical shift and dipolar coupling from their corresponding target SSNMR observables, respectively. It should be noted that the criteria are within the SSNMR experimental errors [21,22]. To generate stable simulation trajectories that often satisfy the aforementioned criteria, the maximum force constants were empirically optimized to be  $k_{cs} = 0.8$  kcal/(mol ppm<sup>2</sup>) and  $k_{dc} = 1.6$  kcal/(mol kHz<sup>2</sup>). However, due to the complexity of the MerF system, we had to slightly lower the criteria ( $\delta\sigma \leq 1.2$  ppm and  $\delta\nu \leq 1.2$  kHz) and increase the force constants ( $k_{cs} = 3.2$  kcal/(mol ppm<sup>2</sup>) and  $k_{dc} = 6.4$  kcal/(mol kHz<sup>2</sup>)) to yield statistically meaningful number of structures.

Fig. 1 illustrates the overall process involved in the structure determination of a monomer using the SSNMR restraint potentials. Starting from an ideal  $\alpha$ -helical conformation ( $\phi = -65^\circ$  and  $\psi = -40^\circ$ ) [23], 130 ps of TAMD were performed 100 times with different initial velocities. In the case of MerF, we generated two parallel TM  $\alpha$ -helices (TM1: residues 26–41 and TM2: residues 51–70) with the helix–helix distance of 15 Å and rotated TM2 from  $-180^\circ$  to  $180^\circ$  in  $5^\circ$  increments. The connecting loop residues were initialized by default CHARMM topologies. We performed 10 TAMD simulations for each initial structure with different initial velocities. The force constants were smoothly increased to the maximum values during a total of 130 ps simulations (150 ps for MerF). The final structures were selected based on the aforementioned criteria.



**Fig. 1.** Illustration of our MD simulation protocol for the SSNMR structure determination. (Right) The conformation changes of fd coat in the course of MD simulations with the SSNMR restraint potentials are shown, starting from an ideal  $\alpha$ -helix (first row) to the final structure (third row) satisfying the RMSD criteria ( $\delta\sigma \leq 1$  ppm and  $\delta\nu \leq 1$  kHz). (Left) The corresponding PISEMA spectra of fd coat are shown. The IP helix is shown in blue and the TM helix in red for both PISEMA spectra and corresponding structures. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

Determination of oligomer structures of VpuTM and M2TMP were initiated from the average structures of the monomers. By displacing the monomer 15 Å away from the origin in a XY direction and rotating it along the Z-axis every  $5^\circ$ , we generated 72 initial structures and ran 10 simulations for each structure with different initial velocities, resulting in 720 assembled structures [24]. By imposing fourfold or fivefold symmetry using the IMAGE facility in CHARMM, 20 ps TAMD simulations were performed with the maximum force constants used in the monomer structure determination. We selected the structures satisfying  $\delta\sigma \leq 1.0$  ppm and  $\delta\nu \leq 1.0$  kHz, and then clustered them based on the helix–helix distance between neighboring (image) helices, the helix–helix crossing angles, and averaged potential energies over last 10 ps. We used the definitions of Chothia et al. [25] for the helix–helix distance and crossing angle.

Solvent/membrane environments play an important role in protein structure. As such, we determined the structures in vacuum as well as in an implicit membrane model in which the average influence of the membrane (waters and lipids) is approximated by solvation free energy terms. As an implicit membrane model, we used a generalized Born model with a simple switching function (GBSW)

[26,27] with 0.04 kcal/(mol Å<sup>2</sup>) for the surface tension coefficient, 23 Å for the hydrophobic thickness corresponding to a DMPC membrane, and 5 Å for a membrane smoothing length over which the hydrophobic region is gradually changed to the solvent region. The implicit planar membrane is perpendicular to the Z-axis and centered at Z = 0. We used a time-step of 1 fs for TAMD simulations with the all-atom parameter set PARAM22 for proteins [28] including the dihedral cross-term corrections (CMAP) [29]. The cutoff distance for non-bonded interactions was 16.0 Å in the monomer structure determination, and there was no cutoff in the oligomer structure determination. A constant temperature of 300 K was maintained by using a simple Berendsen thermostat [30] in combination with velocity reassignment. While we used natural N- and C-terminal groups for fd coat, neutral termini were used for other proteins: acetylated N- and amidated C-terminal groups for M2TMP and MerF and acetylated N- and methylamide C-terminal groups for VpuTM.

## 4. Results and discussion

### 4.1. Fd coat

Fd coat is the major pVIII coat protein of fd filamentous bacteriophage and consists of two distinct  $\alpha$ -helices in the membrane-bound form: one is a short amphipathic in-plane (IP) helix (residues 7–18), which rests on the membrane-water interface, and the other is a longer hydrophobic TM helix (residues 21–45). The structure was determined by solution NMR (PDB:1FDM) [31] as well as SSNMR (PDB:1MZT) [12]. Following the structure determination protocol, we obtained 45 (vacuum) and 61 (GBSW) structures that satisfy the selection criteria ( $\delta\sigma \leq 1.0$  ppm and  $\delta\nu \leq 1.0$  kHz). SI Figs. S4 (vacuum) and S5 (GBSW) show the individual deviations between the calculated SSNMR values and the corresponding experimental data for the best and worst structures satisfying the selection criteria. As shown in Table 2, the fluctuations in calculated tilt angles suggest that the IP helix is more flexible than the TM helix (see the definition of each tilt angle in Table 2). The flexibility of the IP helix may be associated with larger fluctuations in backbone torsion angles of the IP helix (see SI Fig. S6 for

vacuum and S7 for GBSW). The averaged tilt angles of the IP and TM helices (90.1° for IP and 22.4° for TM in vacuum, and 85.3° for IP and 22.0° for TM in GBSW) are in good agreement with those in PDB:1MZT (92.0° for IP and 20.3° for TM) [12]. Fig. 2a shows both the averaged structure obtained in vacuum and PDB:1MZT. A slight bending near residue Ile37 observed in PDB:1MZT is well reproduced. In the MC approach, since the tilt angles of the IP and TM helices are independently determined without the connecting turn (residues 19–20), the resulting structures may have geometric ambiguity. Therefore, Opella and coworkers first categorized their resulting structures into four families that have the same tilt angles but different orientations [12]. They then selected family A (as shown in Fig. 4A of the reference [12]) based on the favorable connecting turn geometry (T19  $\phi/\psi = -93^\circ/-82^\circ$ ; E20  $\phi/\psi = -37^\circ/132^\circ$ ) in the Ramachandran map. Interestingly, 98% of the determined structures (44 of 45) in vacuum are categorized into family A with a similar connecting turn geometry (T19  $\phi/\psi = -112 \pm 12^\circ/-57 \pm 30^\circ$ ; E20  $\phi/\psi = -47 \pm 12^\circ/125 \pm 28^\circ$ ). In GBSW, 82% of the determined structures (50 of 61) are classified into family A. This suggests that our MD approach with SSNMR restraint potentials is more efficient than the MC approach because the former does not require the additional selection process based on the turn geometry.

### 4.2. MerF

MerF is a mercuric ion transporter associated with a bacterial mercury detoxification system. It is an attractive target for the SSNMR structure determination due to its essential biological function of transporting Hg(II) across membranes. MerF is composed of two TM helices (TM1: residues 26–41, TM2: residues 51–70) and an interhelical loop (residues 42–50). The dipolar waves show that MerF has two TM helices with a kink or bend [13]. Table 2 summarizes the features of the structures determined by the simulation protocol and selection criteria. The average tilt angles of the determined structures (TM1/TM2 = 17.3°/18.6° in vacuum and 17.2°/19.8° in GBSW) are in good agreement

**Table 2**  
Structure information of determined monomers

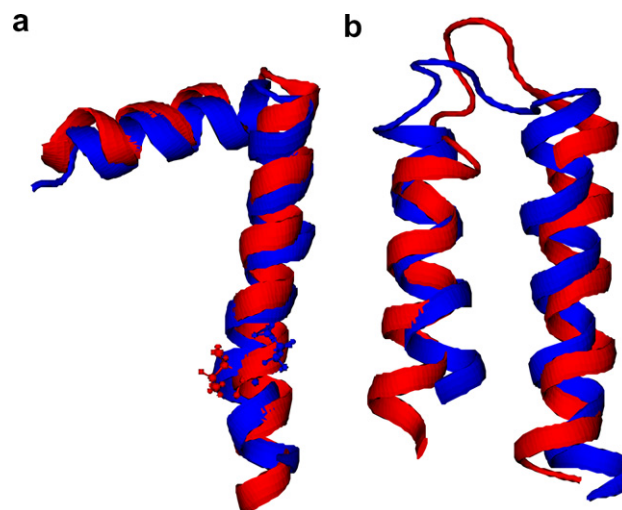
	Fd coat	M2TMP	VpuTM	MerF
<i>Vacuum</i>				
No. of struct. <sup>a</sup>	45	85	21	32
Tilt (TM) <sup>b</sup> (°)	22.4 ± 0.4	31.9 ± 1.0	15.1 ± 0.4	17.3 ± 0.2
Tilt (IP) <sup>b</sup> (°)	90.1 ± 5.0			18.6 ± 0.8
Rotation <sup>c</sup> (°)		173.9 ± 9.6	-27.0 ± 2.6	
Hinge <sup>d</sup> (°)	79.8 ± 8.6			19.4 ± 10.8
<i>GBSW</i>				
No. of struct.	61	100	21	25
Tilt (TM) (°)	22.0 ± 0.4	31.2 ± 0.6	15.0 ± 0.2	17.2 ± 0.4
Tilt (IP) (°)	85.3 ± 5.6			19.8 ± 0.4
Rotation (°)		195.5 ± 6.8	-58.4 ± 2.0	
Hinge (°)	88.8 ± 8.6			19.2 ± 23.0

<sup>a</sup> The number of structures satisfying  $\delta\sigma \leq 1.0$  ppm and  $\delta\nu \leq 1.0$  kHz. In the case of MerF, we used  $\delta\sigma \leq 1.2$  ppm and  $\delta\nu \leq 1.2$  kHz with an additional criterion for the helix-helix distance to be less than 11 Å.

<sup>b</sup> The tilt angle is defined by the angle between the helical principal axis and Z-axis. In the case of MerF, TM corresponds to TM1 and IP to TM2.

<sup>c</sup> People have used two different definitions for the rotation angle: one ( $\rho$ ) from Killian and coworkers [51] and the other ( $\rho_c$ ) from Cross and coworkers [41]. The relationship between the two definitions is  $\rho_c + \rho = -180^\circ$  if  $\rho_c < 0$ . We used the definition of Killian and coworkers with C<sub>α</sub> atoms of Leu26 for M2TMP and Leu11 for VpuTM as reference atoms, respectively.

<sup>d</sup> The hinge angle is defined by the angle between the TM (TM1) and IP (TM2) domains.



**Fig. 2.** Comparison of (a) fd coat and (b) MerF structures (red) determined by our simulation protocol with corresponding structures (blue) previously determined based on the MC conformation search (PDB:1MZT for fd coat and PDB:2H30 for MerF). The fd coat structure is the average of the ensemble structures in vacuum, and the MerF structure is the lowest energy structure that satisfies the additional pre-Pro and Pro  $\phi/\psi$  criteria (see main text). For fd coat structures, Ile37 where the TM helix bends slightly is shown in ball and stick models. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

with those ( $18^\circ$  for both helices) obtained by Opella and coworkers [13]. However, as shown in SI Fig. S8, the determined structures with the present protocol reveal various orientations between two TM helices due to the inherent degeneracy of SSNMR observables as well as the lack of restraints to define their relative orientation. This variation may be associated with the larger fluctuations in backbone torsion angles of the connecting loop than those of the TM1 and TM2 helices although the structures well satisfy the SSNMR observables (see SI Figs. S6 and S7). As such, the best representative structure was selected based on the potential energy as well as the backbone dihedral angles of Leu48, pre-Pro ( $-180^\circ < \phi < -20^\circ$ ,  $30^\circ < \psi < 180^\circ$ ), and Pro49 ( $-110^\circ < \phi < -30^\circ$ ,  $-70^\circ < \psi < 180^\circ$ ) because of their limited  $\phi$  and  $\psi$  angle distribution in the Ramachandran plot [32]. Fig. 2b shows the lowest energy structure in vacuum that satisfies the  $\phi/\psi$  criteria, which is compared with PDB:2H3O determined by Opella and coworkers [13]. It should be stressed that PDB:2H3O was determined by two additional distance restraints between TM1 and TM2 in addition to the SSNMR restraints [13]; our structures were determined by the SSNMR restraints alone based on the aforementioned criteria. The number of final structures in GBSW (18 of 25) is larger than that in vacuum (14 of 32), indicating that the implicit solvent model is useful in enriching the loop structures. However, even with the additional restriction in the  $\phi/\psi$  angles, structural ambiguity remains due to the larger number of residues in the connecting loop. Further structure determination with additional restraints from spectroscopic measurements or with advanced sampling techniques such as the replica exchange method [33,34] may be helpful to obtain the protein structure with higher precision.

#### 4.3. M2TMP

Influenza A virus was the most deadly disease of the 20th century, and its structure determination is still of biomedical importance [35,36]. The TM domain (residues 22–46) of the M2 protein (M2TMP) from influenza A virus forms a proton selective ion channel (homo-tetramer) that is essential for virus infection. The structure of the M2TMP tetramer in bilayers has been studied by experiments and MD simulations [22,24,36–40]. Following the simulation protocol and selection criteria, we first calculated the monomer structure with 10 chemical shifts and seven dipolar couplings (see Table 1). Table 2 summarizes the determined structural features such as TM helix tilt and rotation angles (see the definition of each angle in Table 2). Cross and coworkers estimated a tilt angle of  $37^\circ$  and a rotation angle of  $-10^\circ$  ( $-170^\circ$  or  $190^\circ$  in our definition) in a DMPC bilayer [41]. By using site-directed infrared dichroism spectra, Kukol et al. measured the tilt angle of  $31.6 \pm 6.2^\circ$  [37]. Although the average rotation angle of our structural ensemble is similar to that of the structure determined by Cross and coworkers, the average tilt angle appears to be closer to the measurement by Kukol et al.

Following the simulation protocol for the oligomer structure determination, we obtained 546 (vacuum) and 420 (GBSW) tetramers that satisfy  $\delta\sigma \leq 1.0$  ppm and  $\delta\nu \leq 1.0$  kHz, from which we selected 291 (vacuum) and 145 (GBSW) conformations with a helix–helix distance below  $10 \text{ \AA}$ . The structural data are summarized in Table 3 (vacuum) and in Table 4 (GBSW). As shown in Fig. 3, the low energy structures have left-handed packing (i.e., positive crossing angle) in vacuum as well as in GBSW. The left-handed structure (PDB:1NYJ) determined by Cross and coworkers has Ser31, His37, and Trp41 inside the pore, an important feature supported by the rotational echo double resonance (REDOR) SSNMR spectroscopy [14]. Fig. 4 shows the lowest energy structure in vacuum, that also satisfies the REDOR SSNMR data (i.e., the distance

between His37<sup>(i)</sup> in *i*th helix and Trp41<sup>(i+1)</sup> in the neighboring (*i* + 1)th helix below  $3.9 \text{ \AA}$ ).

Based on the  $\chi_1$  and  $\chi_2$  dihedral angles of His37 and Trp41, the M2TMP tetramer can be categorized into four conformations such as (t60,t90), (t-160,t90), (t-160,t-105), (t60,t-105) where the first “t” represents that the  $\chi_1$  angle is  $180^\circ$  (trans), and the following numerical values are the  $\chi_2$  angles of His37 and Trp41. These conformations yields different distances between His37<sup>(i)</sup> and Trp41<sup>(i+1)</sup> as well as between Trp41<sup>(i)</sup> and Trp41<sup>(i+1)</sup>, which dic-

**Table 3**  
Structure information of determined oligomers in vacuum

	VpuTM	VpuTM	M2TMP
Oligomerization	Tetramer	Pentamer	Tetramer
No. of structures <sup>a</sup>	58	27	291
<i>Population</i>			
Left	28 (48%)	2 (8%)	183 (63%)
Right	14 (24%)	13 (48%)	108 (37%)
Parallel	16 (28%)	12 (44%)	
<i>Crossing angle<sup>b</sup></i>			
Left	$20.1 \pm 0.4$	$16.4 \pm 0.2$	$44.4 \pm 2.0$
Right	$-13.2 \pm 3.0$	$-10.5 \pm 0.8$	$-45.1 \pm 2.2$
Parallel	$8.3 \pm 4.4$	$3.6 \pm 1.2$	
<i>Potential energy<sup>c</sup></i>			
Left	$227.1 \pm 11.6$	$230.0 \pm 6.4$	$-60.9 \pm 10.8$
Right	$244.3 \pm 10.4$	$235.1 \pm 4.6$	$-36.5 \pm 3.6$
Parallel	$277.7 \pm 22.0$	$255.4 \pm 18.4$	
<i>Tilt angle</i>			
Left	$15.0 \pm 0.4$	$15.2 \pm 0.2$	$32.3 \pm 1.6$
Right	$15.0 \pm 0.4$	$15.0 \pm 0.4$	$33.5 \pm 0.8$
Parallel	$15.0 \pm 0.4$	$15.0 \pm 0.2$	
<i>His37–Trp41 distance<sup>d</sup></i>			
Left			$6.1 \pm 1.6$
Right			$20.6 \pm 0.8$

<sup>a</sup> The number of structures satisfying the criteria of  $\delta\sigma \leq 1.0$  ppm and  $\delta\nu \leq 1.0$  kHz, and the helix–helix distance between two neighboring helices below  $10.0 \text{ \AA}$ .

<sup>b</sup> The helix–helix crossing angle was calculated based on the definition of Chothia et al. [25].

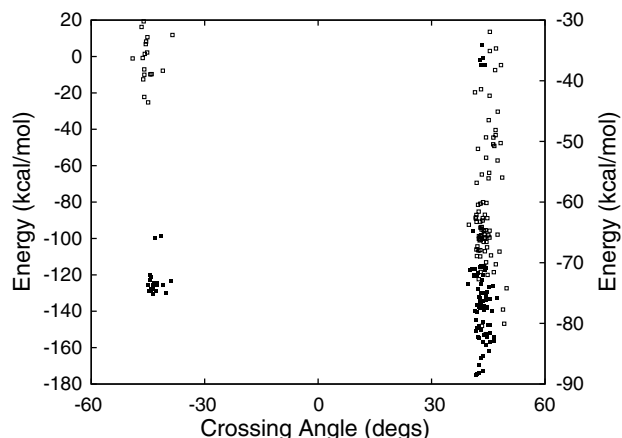
<sup>c</sup> The average potential energy during the final 10 ps simulation.

<sup>d</sup> The distance between N<sub>H</sub>-His37 and C<sub>γ</sub>-Trp41 in M2TMP tetramers.

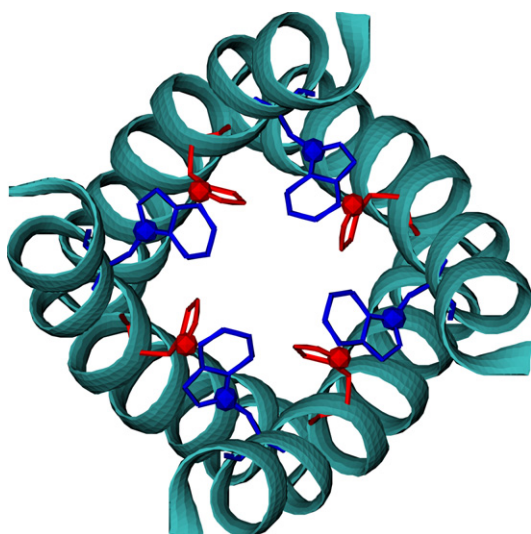
**Table 4**  
Structure information of determined oligomers in GBSW<sup>a</sup>

	VpuTM	VpuTM	M2TMP
Oligomerization	Tetramer	Pentamer	Tetramer
No. of structures	95	101	145
<i>Population</i>			
Left	30 (32%)	11 (11%)	40 (28%)
Right	25 (26%)	26 (26%)	105 (72%)
Parallel	40 (42%)	63 (63%)	
<i>Crossing angle</i>			
Left	$20.4 \pm 0.6$	$16.8 \pm 0.4$	$43.4 \pm 1.4$
Right	$-15.8 \pm 4.2$	$-12.5 \pm 3.2$	$-43.1 \pm 1.6$
Parallel	$2.0 \pm 5.0$	$0.3 \pm 3.8$	
<i>Potential energy</i>			
Left	$195.3 \pm 11.8$	$196.8 \pm 13.4$	$-130.3 \pm 37.2$
Right	$199.7 \pm 7.6$	$206.3 \pm 16.0$	$-123.4 \pm 8.4$
Parallel	$213.8 \pm 13.8$	$221.8 \pm 20.6$	
<i>Tilt angle</i>			
Left	$15.2 \pm 0.4$	$15.3 \pm 0.4$	$31.6 \pm 1.0$
Right	$15.1 \pm 0.4$	$15.3 \pm 0.4$	$31.6 \pm 0.8$
Parallel	$15.1 \pm 0.4$	$15.1 \pm 0.4$	
<i>His37–Trp41 distance</i>			
Left			$8.3 \pm 2.8$
Right			$19.4 \pm 0.6$

<sup>a</sup> The same definitions as those in Table 3 are used for all the quantities listed.



**Fig. 3.** Distributions of helix-helix crossing angles and potential energies of M2TMP tetramer structures in vacuum (empty squares, right Y-axis) and GBSW (filled squares, left Y-axis). Only first 100 low energy structures are shown for clarity.



**Fig. 4.** A M2TMP tetramer structure with the lowest potential energy in vacuum, satisfying  $\delta\sigma \leq 1.0$  ppm and  $\delta\nu \leq 1.0$  kHz as well as the REDOR SSNMR data. The distance between  $N_\alpha$  – His37 (red sphere) and  $C_\gamma$  – Trp41 (blue sphere) is 3.62 Å. The helices are shown in cyan. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

tates the relative orientations of both residues inside the pore; for example, based on PDB:1NYJ, these distances are (3.8 Å, 11.7 Å) in (t60,t90), (3.6 Å, 11.4 Å) in (t-160,t90), (3.8 Å, 5.1 Å) in (t-160, t-105), and (3.6 Å, 5.1 Å) in (t60,t-105), where the first one is for His37<sup>(i)</sup> and Trp41<sup>(i+1)</sup> and the second one for Trp41<sup>(i)</sup> and Trp41<sup>(i+1)</sup>. Three conformations corresponding to (t-160,t90), (t-160,t-105), and (t60,t90) have been determined by <sup>19</sup>F SSNMR [42], REDOR SSNMR [14] and MD simulations [43], respectively.

In our calculations, when the initial  $\chi_2$  of Trp41 was set to 0°, which is the default CHARMM topology, the determined structures were populated mostly in (t60,t-105) and partly in (t-160,t-105) conformations. However, when the  $\chi_2$  was set to 90°, the determined structures were mainly in (t60,t90) and partly in (t-160,t90) conformations. Recently, the DeGrado group determined the X-ray structure of the M2TMP tetramer (PDB:3BKD) [6]. The structure forms left-handed helices with tilt angles of 35° ± 2°. However, apart from the asymmetric nature of the tetramer structure, the average distance between His37<sup>(i)</sup> and Trp41<sup>(i+1)</sup> is 13.4 ± 1.0 Å which does not agree with most SSNMR experiments.

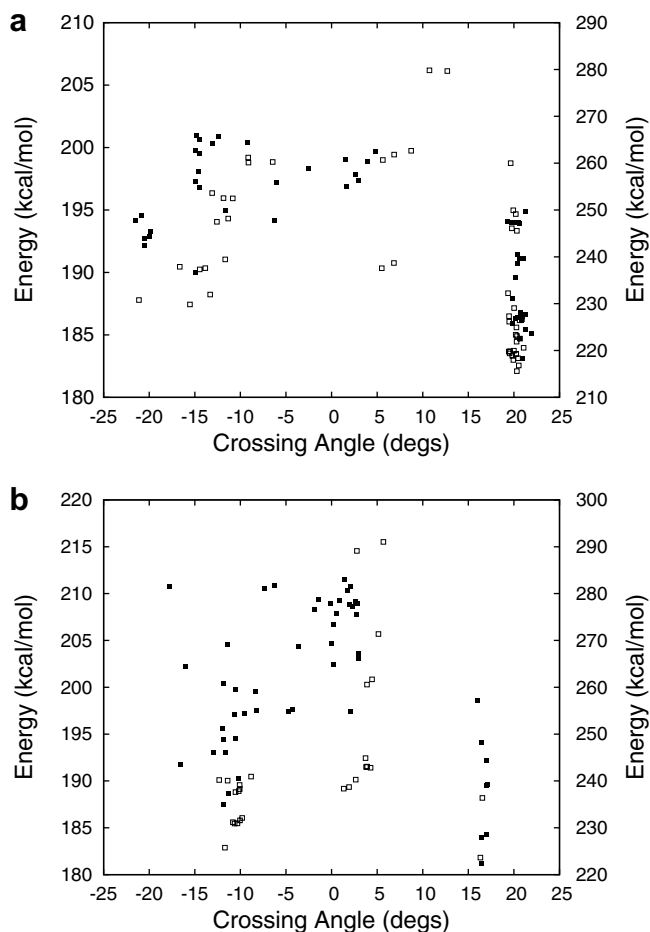
In addition, the X-ray structure has two  $\chi_1$  and  $\chi_2$  conformations such as (t-160,t90) and (t60,t90). Clearly, further experimental and computational studies will be required to understand the difference of the interfacial helix packing and the orientations of pore-lining residues such as His37 and Trp41.

#### 4.4. VpuTM

Viral protein “u” (Vpu) is a small membrane protein whose sequence is encoded in the genome of HIV-1. There is a hydrophobic TM helix in the N-terminal domain and amphipathic helices in the C-terminal domain [44]. Vpu affects the budding of new virus particles, and the hydrophobic TM helix (residues 7–25) may play an important role as an ion channel that is selective for monovalent cations such as Na<sup>+</sup> and K<sup>+</sup> [45]. The ion channel is formed by the homo-oligomers of four to seven helices. Computational studies have suggested that the most probable oligomeric state is a pentamer [46,47]. However, a recent study suggested that there might be more than one possible oligomeric state [9], and the exact state is yet to be determined. The atomic structure of Vpu is an essential and crucial starting point for rational drug design [48]. Vpu has also served as a model system (VpuTM) for method development in structure determination of membrane proteins [9,49]. As summarized in Table 2, there are 21 determined monomer structures in both vacuum and GBSW membrane. The average tilt angles in both vacuum and GBSW are around 15°, which agrees well with 13° determined by the Opella group [9]. Interestingly, the fluctuations of individual backbone torsion angles of VpuTM are smaller than those of M2TMP probably due to a larger number of SSNMR observables used in VpuTM than M2TMP (see SI Fig. S6 for vacuum and S7 for GBSW).

By using our simulation protocol for oligomer structure determination, we obtained 58 tetramers and 27 pentamers in vacuum, and 95 tetramers and 101 pentamers in GBSW membrane model. The distribution of averaged energies as a function of crossing angle is shown in Fig. 5a (tetramer) and b (pentamer). While M2TMP tetramers are classified into two conformations (left- and right-handed), VpuTM oligomers are categorized into three conformations: left-handed ( $\Omega > 15^\circ$ ), right-handed ( $\Omega < -7.5^\circ$ ), and parallel ( $-7.5^\circ \leq \Omega \leq 15^\circ$ ), where  $\Omega$  is the crossing angle between two neighboring helices. The structural analyses of these conformations are summarized in Table 3 (vacuum) and Table 4 (GBSW). The parallel structures are energetically unfavorable for both tetramer and pentamer, as shown in Fig. 5a and b. The left-handed conformations in tetramer and pentamer are energetically more favorable than the right-handed conformation in both vacuum and GBSW, although the energy difference is small. In the left-handed conformation, Trp22 is located inside the pore and Ser23 outside the pore. Fig. 6 shows three pentameric conformations obtained in GBSW.

Based on rigid-body assembly calculations with the AMBER potential, the Opella group proposed that both oligomers (tetramer and pentamer) have the strong preference of being the right-handed conformation due to energetically unfavorable clashes of the large Trp22 sidechains inside the pore in the left-handed conformation [9]. They also argued that the resulting pore lining of Ile17 in the right-handed conformation can be supported by the fact that the highly conserved Ile17 may have an essential role in its ion channel activity. However, modeling based on the experimental data from site-directed Fourier transform infrared dichroism supports the left-handed pentamer conformation and reveals a pore occluded by Trp residues at the end of the TM domain [50]. Our assembly calculations with flexible sidechains show that the left-handed pentamer conformation can have a slightly lower energy structure than the right-handed conformation without clashes of the Trp22 sidechains inside the pore. Clearly, further



**Fig. 5.** Distributions of helix–helix crossing angles and potential energies of VpuTM (a) tetramer and (b) pentamer structures in vacuum (empty squares, right Y-axis) and GBSW (filled squares, left Y-axis). Only first 50 low energy structures (27 pentamer structures in vacuum) are shown for clarity.

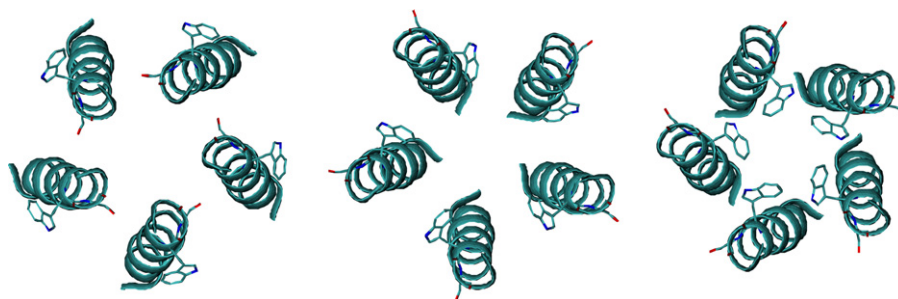
experimental measurements are required to resolve this discrepancy.

## 5. Conclusions

We have formulated and implemented SSNMR restraint potentials for  $^{15}\text{N}$  chemical shift and  $^{15}\text{N}$ – $^1\text{H}$  dipolar coupling. A series of assessments shows that the calculated restraint forces are numerically accurate and stable. The simulation stability depends on restraint force constants as well as the simulation protocol. Combined with constant temperature TAMD simulations, we have

also developed a simple SSNMR structure determination protocol with these restraint potentials. The protocol consists of (1) generating an ideal  $\alpha$ -helix (or helices), (2) performing several steps of TAMD simulations with increasing force constants, and (3) selecting structures that satisfy certain criteria such as RMSD between calculated and target SSNMR observables, i.e.,  $\delta\sigma \leq 1.0$  ppm and  $\delta\nu \leq 1.0$  kHz. Its efficacy in determining the relative orientations of TM helices with respect to membrane bilayers has been illustrated by previously characterized four membrane proteins: (1) fd coat protein from fd filamentous bacteriophage, (2) MerF from bacterial mercury detoxification system, (3) M2TMP from Influenza A virus, and (4) Vpu TM domain from HIV-1. For M2TMP and VpuTM, we further demonstrated that the restraint potentials can be combined with the IMAGE facility in CHARMM to model corresponding homo-oligomer structures: tetramer for M2TMP, and tetramer and pentamer for VpuTM.

The structures determined by our protocol are characterized by tilt and rotational angles of the TM helices and show overall good agreement with the published structures determined using the existing MC approach. Here, we briefly summarize the main findings from our structure calculations. In the case of fd coat, when the determined structures were clustered, we found less geometric ambiguities in the relative orientation between the IP and TM helices despite the degeneracy inherent in orientational restraints. This might be attributed to the potential energy function which favors closer interactions between the IP and TM helices. The determined MerF monomer structures show structural degeneracy in the relative orientation between two TM helices, although the tilt and rotational angles of individual TM helices are consistent with previously determined structures. Additional restraints from more spectroscopic measurements appear to be necessary to resolve the degeneracies. Identifying proper monomer structures greatly reduces the conformational space needed to be searched for corresponding oligomer structures. By clustering the determined oligomers based on total potential energies and crossing angles, we found that the M2TMP tetramer has strong preference of being left-handed, placing Ser31, His37, and Trp41 in the pore region, which is in good agreement with the previously determined structure [14,42] and the recent X-ray structure [6]. However, based on the  $\chi_2$  distributions of His37 and Trp41, four different conformations were identified. Further studies are needed to identify the most probable conformation or energetic difference in four conformations. Although the oligomer structure of VpuTM has not been determined yet at the atomic level, we proposed that the favorable conformations of the VpuTM tetramer and pentamer are left-handed structures. This is consistent with the Fourier transform infrared dichroism experiment, but is at odds with the structure previously proposed based on rigid-body conformational searches using available SSNMR observables and the AMBER force field. Additional restraints from other spectroscopic measurements may help to resolve the discrepancy.



**Fig. 6.** Three conformations of VpuTM pentamers with the lowest potential energy in right-handed, parallel, and left-handed structures (from left to right) in GBSW. The helices are shown in cyan, and Trp22 and Ser23 residues are shown in stick models.

Finally, we also examined the influence of incorporating solvation effects in our structural calculation. The results show that the number of structures satisfying the selection criteria is generally larger in structure calculations with the GBSW implicit membrane model than in vacuum. This indicates that structure calculations with implicit solvent models might be useful in structure determination and modeling of membrane proteins with a limited number of available experimental observables. Further computational experiments with advanced sampling techniques such as the replica exchange method [33,34] might provide better insights on how these techniques can be combined to yield biological useful models using limited number of orientational restraints.

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## Appendix A. Supplementary data

Supplementary data associated with this article can be found, in the online version, at [doi:10.1016/j.jmr.2008.04.023](https://doi.org/10.1016/j.jmr.2008.04.023).

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